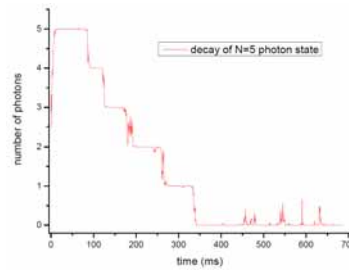





Experimental realizations of quantum trajectories: the quantum jumps of light

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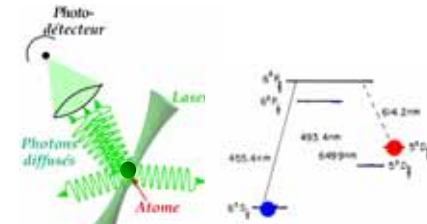
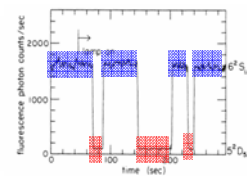
Quantum jumps

- A striking quantum trajectory
 - a single quantum system
 - slowly evolving
 - frequently measured by an ideal projective quantum measurement
 - Evolves, in a single realization of the experiment, by sudden jumps between eigenstates of the measured observable
- Evolution in a particular trajectory
 - differs from classical evolution
 - differs from usual continuous evolution predicted by quantum statistical averages.





Quantum jumps for a trapped ion

- A paradigmatic example of quantum jumps



Nagourney et al. (1986)

- First experimental realization of quantum trajectories
- Quantum jumps also observed for other matter particles (electrons, molecules, atoms, artificial qubits)
- In all these experiments, a matter particle is interrogated by light
- Is it possible to observe the quantum jumps of light?
 - need an ideal photodetector, realizing a projective measurement

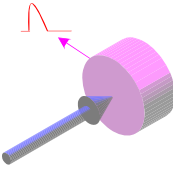
Ideal quantum measurement

- The most intriguing aspect of quantum theory
- Simple postulates for an ideal (projective) measurement
 - Quantum discontinuity
 - not all results allowed
 - » eigenvalues of the measured observable
 - Statistical results
 - predict only measurement results probabilities
 - » ‘God is playing dice’
 - State collapse and repeatability
 - two identical measurements in a short time interval always give identical results
 - » state projected onto an eigenstate of the observable





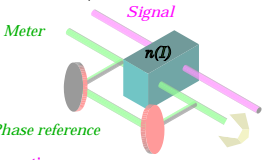
Ideal and real quantum measurements

- Most quantum measurements are far from ideal
 - e.g. Photodetection (counting photons)
 - measurement of light field energy
 - quantized result: number of photons
 - statistical: photon number statistics
 - repeatable?
- Photodetectors (PM's, photodiode, retina) absorb incoming photons, converting their energy into an electrical/chemical signal.
 - A second detection always gives zero: impossible to 'see' the same photon twice
 - The field state is demolished by the detection
- This demolition is not a requirement of quantum mechanics




Ideal photon number counting

- Quantum non-demolition measurements (Braginsky, 70s)
 - A transparent photocounter
 - 'see' the same photon twice
 - provides vision to the invisible man...
 
- Realized in the optical domain (Grangier et al, Nature, 396, 537)
 - no single photon resolution
 - weak non-linearity
 - propagating fields:
 - repetition difficult
 - Not appropriate for quantum jumps observation



Quantum jumps of photons?

- A QND photodetector operating at the individual photon level
- A photon 'box' able to store a photon for a long time
 - back to Einstein's dream



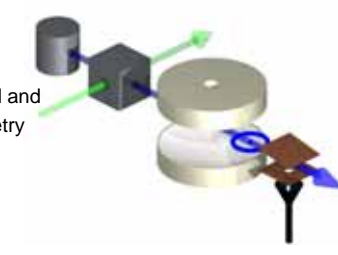
Yet another gedankenexperiment



- A clock whose ticking rate is determined by the number of photons in a box
- The final clock hand's position directly measures the photon number

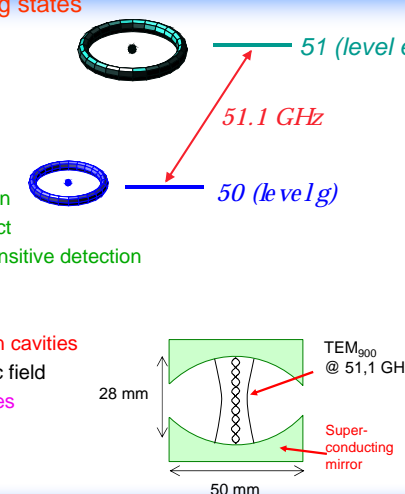
Our boxes and clocks

- Photon box
 - a superconducting microwave cavity
 - longest available photon storage time
- Clock
 - circular Rydberg states
 - state superpositions prepared and probed by atomic interferometry
- Interaction
 - Cavity quantum electrodynamics
 - spin/spring system
 - matter-field coupling at its simplest
 - also in optical CQED, circuit QED, ion trap physics



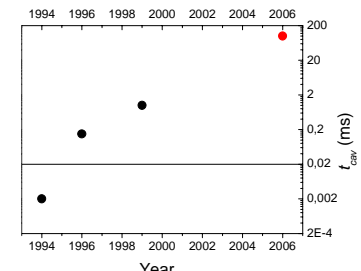
Circular Rydberg states

- Nearly ideal atoms
 - Mesoscopic orbit size
 - 0.25 μm diameter
 - Long lifetime (30ms)
 - Large coupling to radiation
 - tunable via the Stark effect
 - Efficient (> 80%) state sensitive detection
- compatible only with open cavities
 - static directing electric field
- Fabry Perot (open) cavities
 - essential parameter
 - photon lifetime



A box for microwave photons

- optimization of the cavity quality
 - a long (painful !!) process
 - our pet Moore's law



Year	τ_{cav} (ms)
1994	~0.002
1996	~0.1
1998	~1
2006	~100

- extrapolations might not be safe....

A box for microwave photons

- An ultra-high finesse cavity
 - Best lifetime: 0.13 s @ 0.8 K

- quality factor:
 - $Q = \omega T_c = 4.2 \cdot 10^{10}$
- Finesse
 - $F = 4.6 \cdot 10^9$

- during this macroscopic time interval a photon
 - bounces 1.1 billion times on the mirrors
 - travels 40 000 km
 - » if it was in an optical fiber: attenuation 0.00011 dB/km !!
- Best mirrors so far, in any frequency domain !

S. Kuhr et al, APL **90**, 164101 (2007)

A new mirror technology




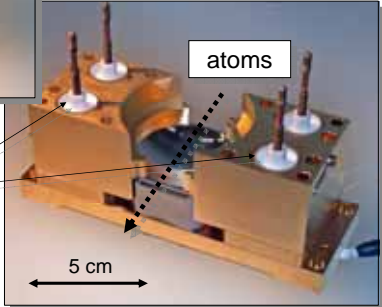
- Copper substrates
diamond machining
~shape accuracy 300 nm pvt
~rugosity 10 nm
Toroidal surface → single mode



- 12 μm Niobium layer
Cathode plasma sputtering
CEA, Saclay
[E. Jacques, B. Visentin, P. Bosland]

S. Kuhr et al, APL, 90, 164101

Cavity assembly

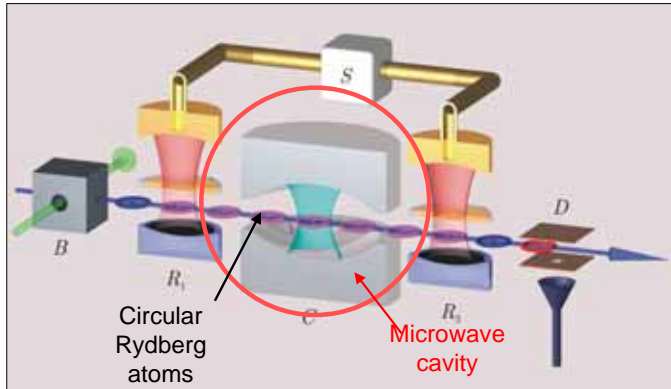



atoms

piezos

5 cm

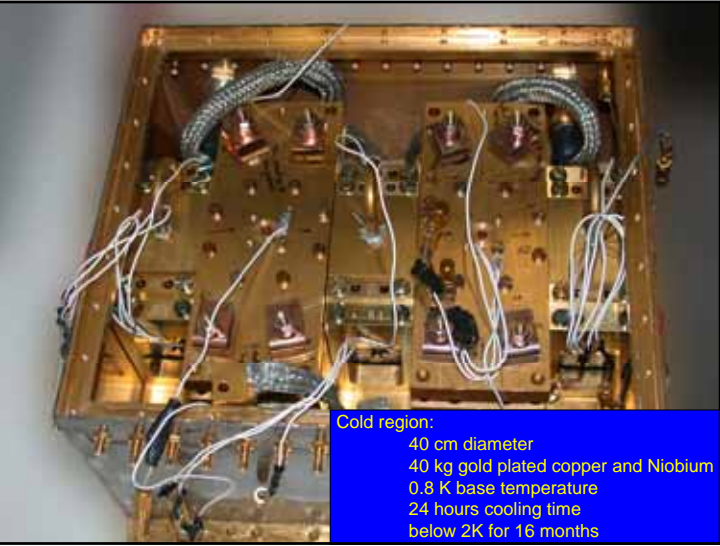
Experimental set-up



Circular Rydberg atoms

Microwave cavity


RMP 73, 565



Cold region:
40 cm diameter
40 kg gold plated copper and Niobium
0.8 K base temperature
24 hours cooling time
below 2K for 16 months

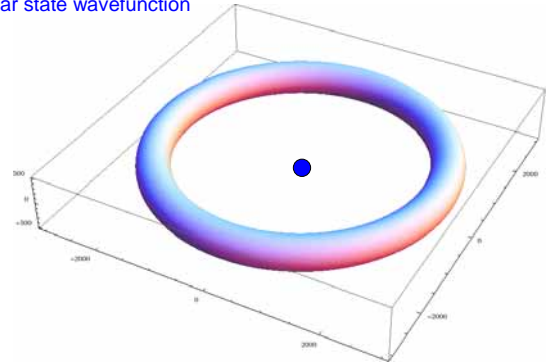
A QND photon measurement process

- How to 'see' a photon with an atom without absorbing it ?
 - non-resonant (dispersive) atom-field interaction
 - no energy exchange
 - no photon emission or absorption
 - Light shifts
 - atomic levels shifted by the photon field
 - Photon number-dependent light shifts
- Find a clock's hand on the atom ?



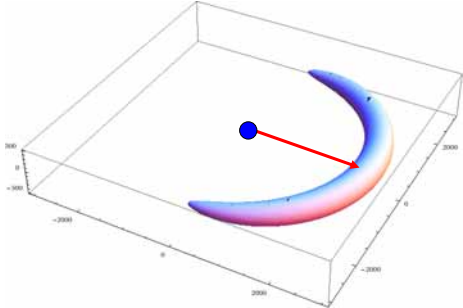
An atomic clock

- Circular state wavefunction
 - Electron delocalized on Bohr's orbit: No phase information
 - Zero atomic dipole
 - Not a very good clock: no hands...



An atomic clock

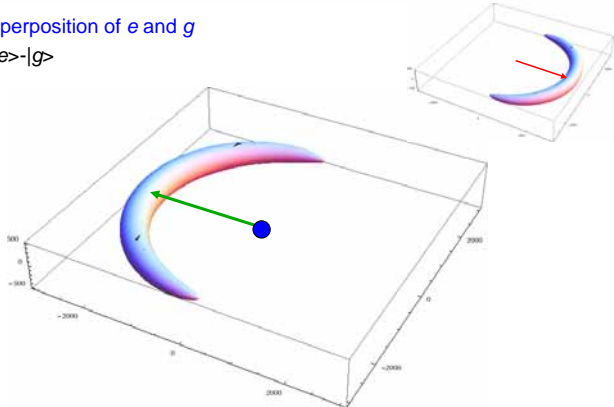
- A superposition of e and g
 - $|e\rangle+|g\rangle$
 - Prepared by a $\pi/2$ pulse of resonant classical microwave



- 'well defined' phase of the electron on the orbit
- Non-zero electric dipole: the clock's hand

An atomic clock

- A superposition of e and g
 - $|e\rangle-|g\rangle$

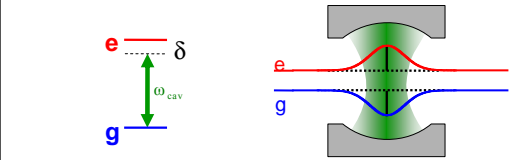


- Opposite dipole
- quantum state orthogonal with $|e\rangle+|g\rangle$

An atomic clock

- Distinguishing orthogonal states
 - A second $\pi/2$ pulse (identical to that preparing $|e\rangle+|g\rangle$) transforms
 - $|e\rangle+|g\rangle$ into $|g\rangle$
 - Addition of the two pulses
 - $|e\rangle-|g\rangle$ into $|e\rangle$
 - Subtraction of the two pulses
 - A consequence of unitarity
 - Different final atomic energies, easily distinguished by the field ionization detection
 - Note: the energy for the transition is provided by the classical resonant microwave field

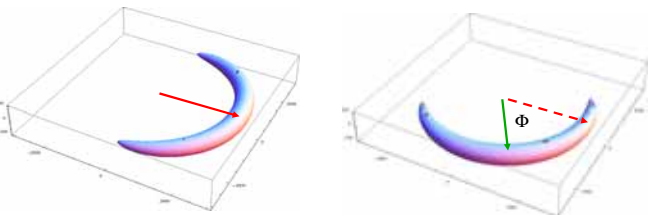
Dispersive atom-field interaction



- Atomic frequency shift inside the cavity
 - 'clock' ticking rate modification $\frac{\Omega_0^2}{2\delta}(n+1/2)$ $\Omega_0/2\pi = 50$ kHz
 - $\frac{1}{2}$ contribution: cavity-induced Lamb shift effect (PRL 72, 3339)
 - fixed shift: no influence on QND measurement
 - n contribution: light shift at the single photon level (also observed in circuit QED -- Nature 445, 515)
 - Adiabatic coupling in and out of the atom-cavity interaction
 - negligible spurious absorption rate ($<10^{-4}$ for $\delta \sim \Omega$)

Phase shift of the atomic dipole

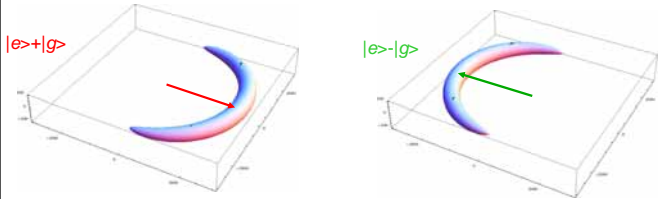
- A $\pi/2$ microwave pulse prepares the $|e\rangle+|g\rangle$ superposition before atom-cavity interaction
- Light shift accumulated over atom-cavity interaction time t_i
 - A phase shift of the atomic dipole

$$\Phi = \Phi_0 n \quad \Phi_0 = \frac{\Omega_0^2}{2\delta} t_i$$


- Final states are generally not orthogonal

A simple case

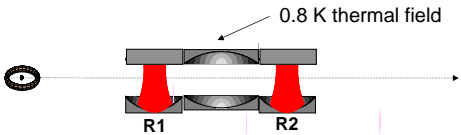
- Field containing zero or one photon
 - π phase shift per photon for the dipole: $\Phi_0 = \pi$
 - Two orthogonal final states



- Zero photon (no shift)
- One photon (π phase shift)
- A second $\pi/2$ pulse of resonant microwave leads to e and g
 - g for zero photon
 - e for one photon

Repeated measurement of a small thermal field

- Probe the cavity equilibrium thermal field
 - 0.8 K, 0.05 photon on the average

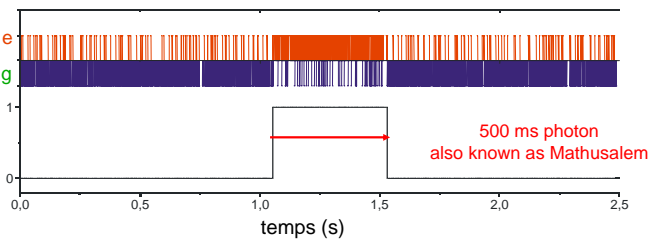


- 50 000 atomic probe samples
 - separated by 70 μ s
 - over a 3.5 s time interval
 - each undergoing the full Ramsey sequence
- 0.1 atom in each sample (no two-atoms events)

Gleyzes et al, Nature, 446, 297 (2007)

Birth, life and death of a photon

$T=0.8$ K $n_{th}=0.05$

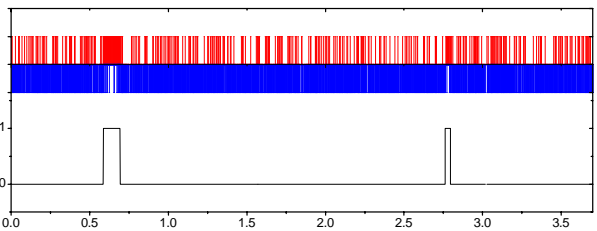


500 ms photon also known as Mathusalem

temps (s)

Gleyzes et al, Nature, 446, 297 (2007)

Birth, life and death of two photons

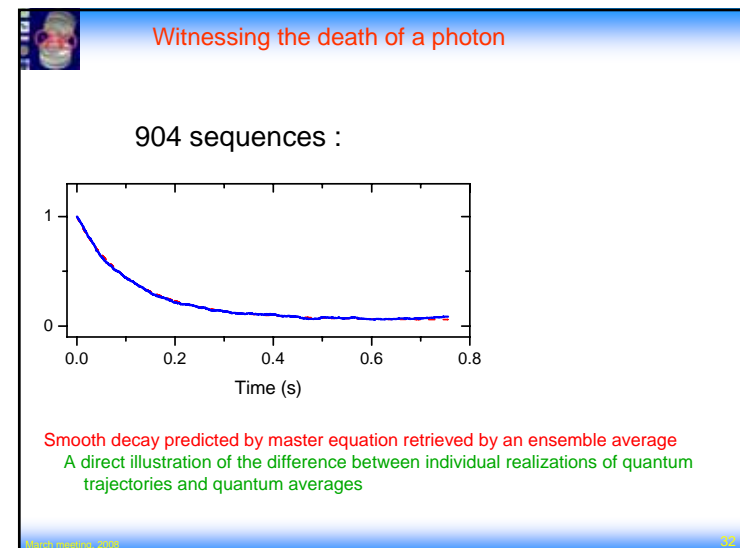
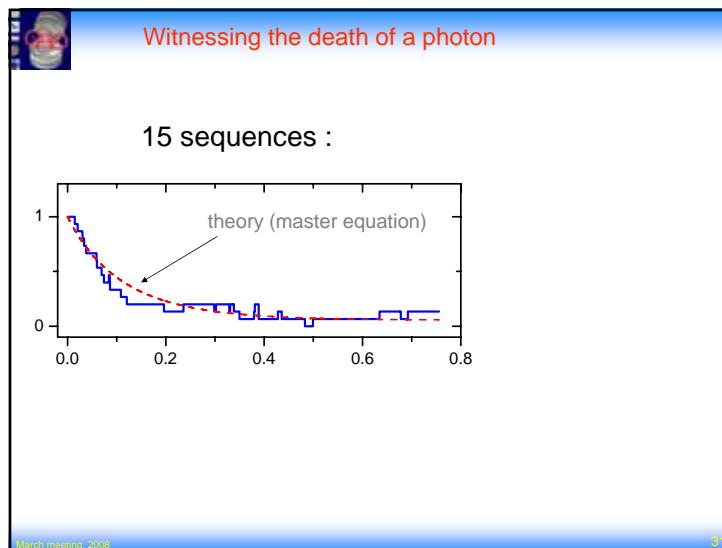
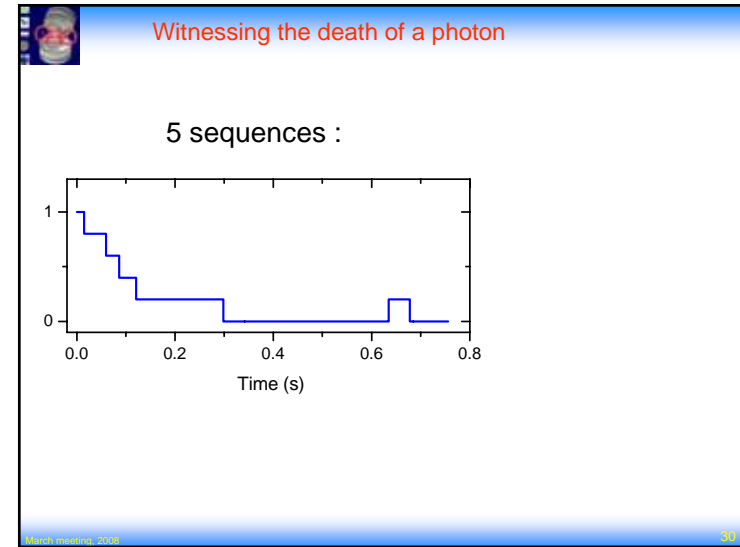
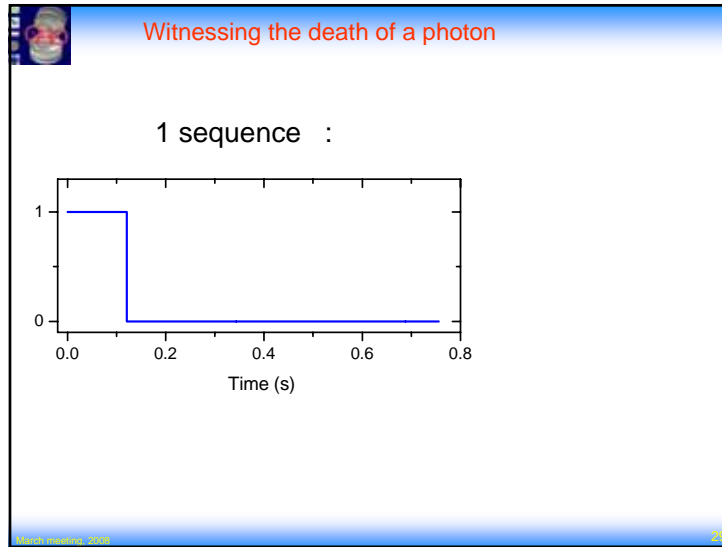


Gleyzes et al, Nature, 446, 297 (2007)

Witnessing the death of a photon

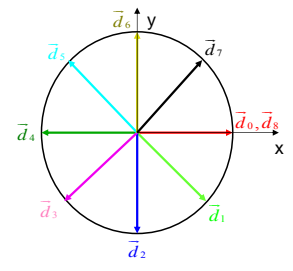
- Quantum trajectories starting from $n=1$:
 - Erase cavity field with atomic samples in g , tuned at resonance (via the Stark effect). High efficiency preparation of the vacuum state.
 - Prepare the one-photon Fock state
 - emission of a single atom prepared in e , tuned at resonance
 - Repeatedly probe the cavity
 - 10 000 atomic samples over 700 ms

Gleyzes et al, Nature, 446, 297 (2007)



Counting up to seven ?

- $\Phi_0 = \pi/4$ phase shift per photon
 - Eight different final orientations of the atomic dipole



Quantum ambiguity: non-orthogonal final states
How can we distinguish them ?

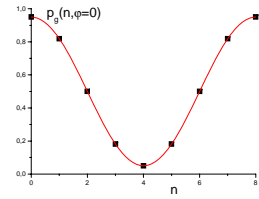
Final detection probabilities

- State after a second $\pi/2$ pulse (identical to the first, phase $\varphi=0$)

Probability $p_g(n, \varphi)$ for detecting g :

$$p_g(n, \varphi=0) = \frac{1}{2} (1 + \bar{d}_n \cdot \bar{u}_x)$$

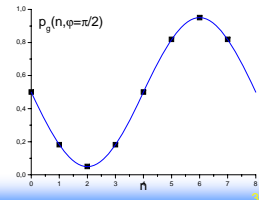
Linear function of the x component of the (normalized) atomic dipole


- State after a second $\pi/2$ pulse (in quadrature with the first, phase $\varphi=\pi/2$)

Probability $p_g(n, \varphi)$ for detecting g :

$$p_g(n, \varphi=\pi/2) = \frac{1}{2} (1 + \bar{d}_n \cdot \bar{u}_y)$$

Linear function of the y component of the (normalized) atomic dipole

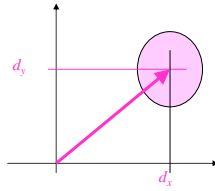


One atom is not enough

- Knowledge of $p_g(\varphi=0)$ and $p_g(\varphi=\pi/2)$ determines d_x and d_y and hence n
- A single atomic detection provides only binary information
 - e or g
 - Not enough information to determine n
- Two no-go theorems
 - The state of a single quantum system cannot be measured
 - A single bit is not enough to count from zero to seven
- More than one atom needed to determine the photon number

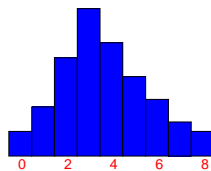
Counting by accumulating information

- Initial cavity state: $\sum_n c_n |n\rangle$
- One-atom cavity state: $\sum_n c_n |n\rangle \otimes |\bar{d}_n\rangle$
- N atoms-cavity state: $\sum_n c_n |n\rangle \otimes (|\bar{d}_n\rangle)^{\otimes N}$
 - Entanglement of the photon number with a mesoscopic atomic sample
 - All atoms have the same dipole orientation for a given photon number
- Split atomic sample in two parts
 - On $N/2$ atoms, second $\pi/2$ pulse with $\varphi=0$
 - Measure p_g i.e. x-component of dipole, d_x
 - On $N/2$ atoms, second $\pi/2$ pulse with $\varphi=\pi/2$
 - Measure p_g i.e. y-component of dipole, d_y
 - Estimate dipole direction with $1/\sqrt{N}$ uncertainty



Estimating the dipole direction

- A measurement sequence
 - initialize the cavity field
 - remove leftovers of previous experiments
 - inject a 3.5 photon coherent field
 - Poisson photon number distribution
 - measure the spin direction
 - send 110 atoms
 - analyze the final state with two Ramsey phase settings
 - compile the measurement results



Dipole directions

- One measurement

Dipole directions

- Thousands of measurements

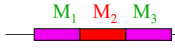
- Rotation of the dipole
 - evidence of light shifts
- Discrete privileged directions:
 - field intensity quantized !!
 - quantum discontinuity
- Probability distribution
 - reveals the photon distribution
 - more on that soon
- Looks like a measurement
 - is it repeatable ???

Sequential measurements

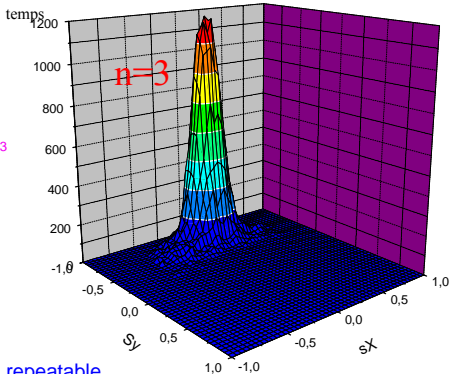
- Three adjacent 110 atoms samples
 - three dipole measurements
 -
 - select those measurements such that $n=3$ for M_1 and M_3 (get rid of cavity relaxation)
 - plot the statistics of dipole directions for M_2

Sequential measurements

- Three adjacent 110 atoms samples



– M2 always give the same result as M1/M3



- The measurement is repeatable
 - ideal projective measurement of photon number !!

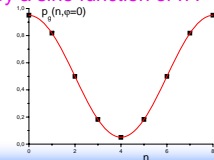
Progressive collapse of the field state

- Another view on the N atoms measurement sequence
 - each atom brings partial information on the field
 - each detection changes our knowledge of the field
 - each detection changes the photon number distribution we can infer
 - » Initial inferred distribution flat (no information)
 - » Each atomic detection changes it
 - » Final distribution (after enough atoms recorded): a single Fock state
- Describe this progressive information acquisition by the Bayes law
 - an information-theory point of view equivalent to the state projection postulate.

C. Guerlin et al, Nature, 448, 889

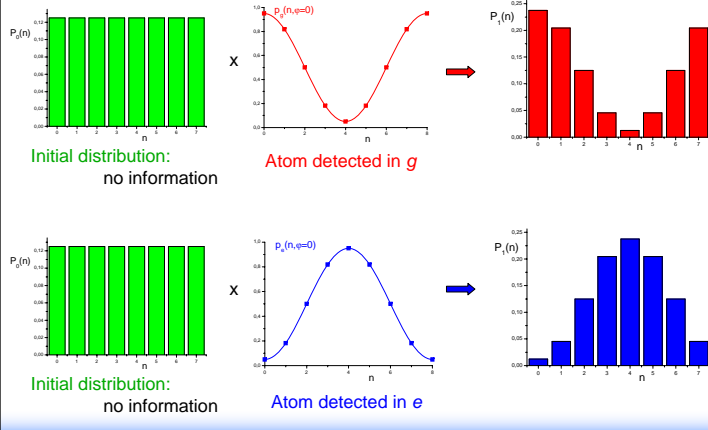
An elementary step in the collapse

- Inferred probability distribution after k atoms have been detected $P_k(n)$
- $k+1$ th atom analyzed with phase φ for second $\pi/2$ pulse
 - Found to be in level j ($j=e$ or g)
- New photon number distribution: $P_{k+1}(n) = P_k(n) \frac{P_j(n, \varphi)}{P_j(\varphi)}$
 - where $P_j(\varphi) = \sum_n P_k(n) P_j(n, \varphi)$ is the *a priori* probability for getting the atom in state j
- At each step, the inferred distribution is multiplied by a sine function of n !
 - Some photon numbers become less likely
 - Photon number distribution decimation



A photon number decimation process

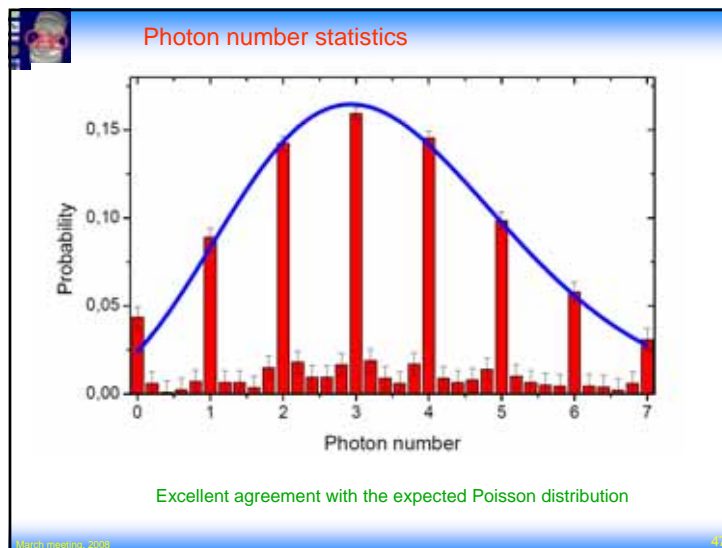
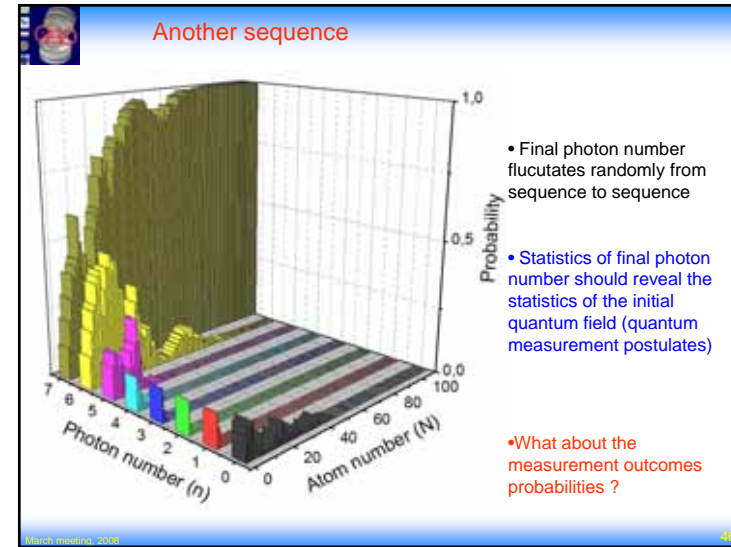
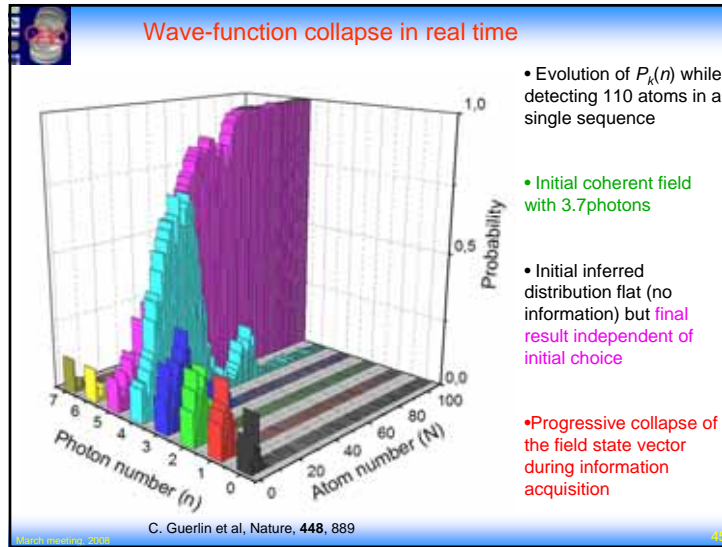
- First step in the decimation process ($k=0$)



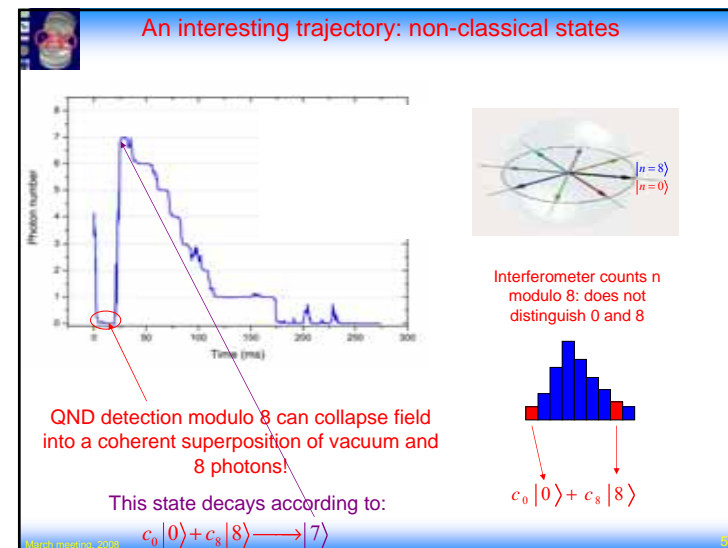
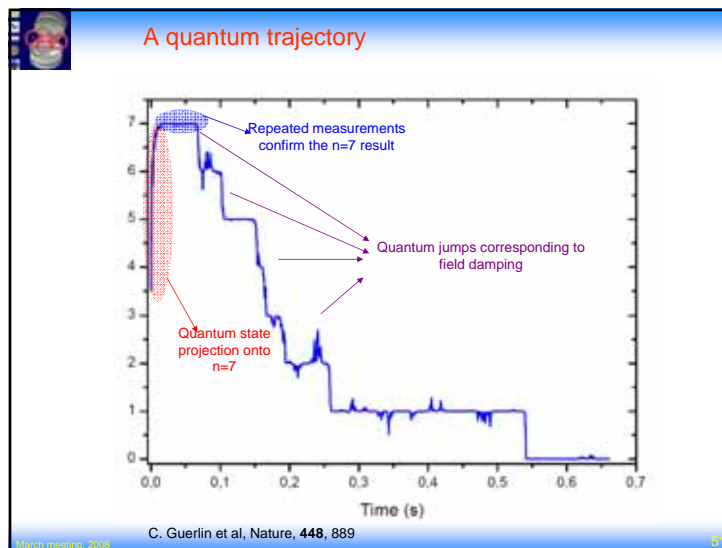
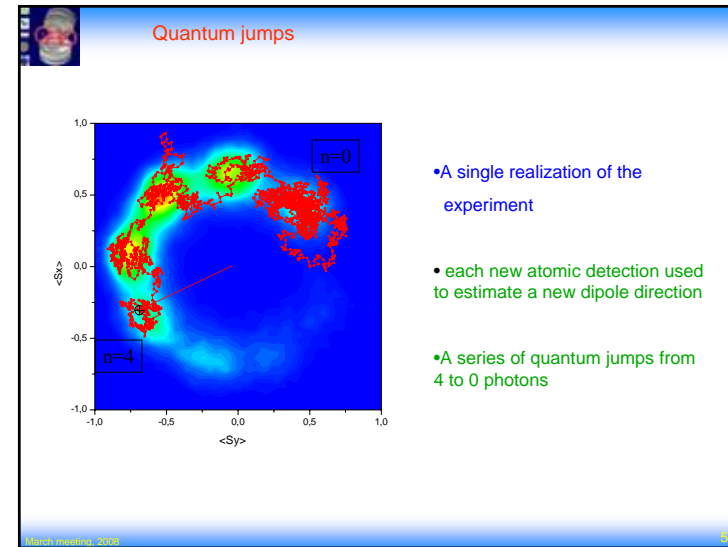
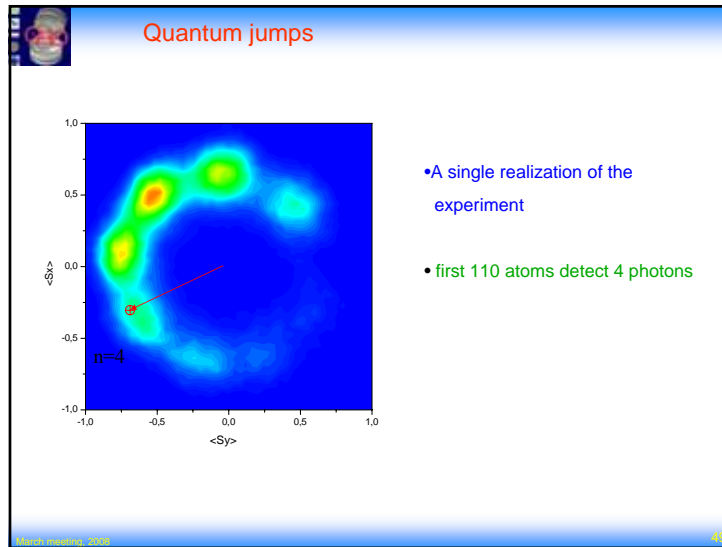
Initial distribution: no information

Atom detected in g

Atom detected in e



- ### Quantum jumps of light
- Keep measuring the photon number over a long time
 - prepare a 3.5 photons coherent field
 - send atoms repeatedly
 - at each time infer the photon number from the last 110 atomic detections
 - dipole direction measurement
 - decimation process



Measurement and dynamics

- For an incoherent (relaxation) process, the QND measurement does not affect the dynamics
 - Relaxation is not sensitive to coherence
 - Lifetime of the cavity states unchanged by the continuous monitoring
- A completely different situation for a coherent evolution

Quantum Zeno effect

- A watched kettle never boils
 - coherent evolution of a system and frequently repeated quantum measurements
 - a quantum jumps evolution between eigenstates of the measured quantity
 - an evolution much slower than without measurements
 - no evolution at all in the limit of zero delay between measurements

e.g. coherent evolution of a two-level system

Note: no Zeno effect for relaxation processes

Quantum Zeno effect

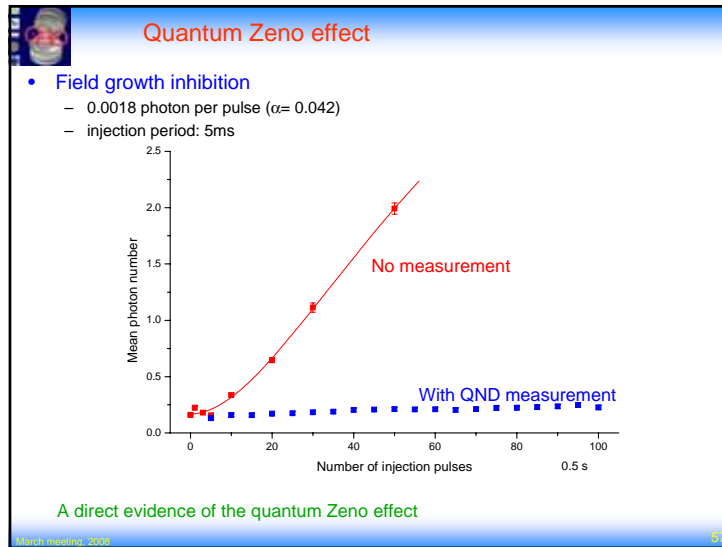
- Coherent evolution: injection of a coherent field by a classical source

– Repeated injection of phase coherent pulses: an amplitude varying linearly with the number of injections (photon number varies quadratically).

Principle of the experiment: perform QND measurements of photon number between two pulses

A difficult experiment

- Phase coherent injections over a second
 - excellent control of cavity frequency (sub Hz stability)
 - The cavity is a quite sensitive multimeter!
 - removal of acoustic noise, vibrations insulation
 - temperature stabilization within 0.1 mK
 - He bath pressure regulation below 1 mb!
 - PZT voltage regulated to 0.1 mV
 - <0.3Hz drift over 15 minutes
 - Measured cavity linewidth matches damping time
 - inferred phase noise of a few degrees only over T_c



- ### Conclusion and perspectives
- Quantum jumps of light observed for the first time
 - a direct illustration of quantum measurement postulates
 - a single quantum controls the states of hundreds of atoms
 - promising for quantum information processing
 - preparation of highly non-classical field states
 - high photon number Fock states
 - Schrödinger-cat like superpositions
 - Perspectives
 - exploration of the non-classical states
 - direct measurement of their Wigner function
 - decoherence studies
 - Towards non-local cat states shared by two cavities
 - interplay of Bell inequalities with decoherence

The ENS team

- Permanent members:
 - S. Haroche, M. Brune, G. Nogues, JM Raimond
- Cavity QED experiments
 - S. Kuhr, I. Dotsenko
 - S. Gleyzes, C. Guerlin, J. Bernu, S. Deléglise, C. Sayrin
- Superconducting atom chip
 - A. Lupascu
 - T. Nirrengarten, C. Roux, A. Emmert
- www.cqed.org
- €€
 - JST (ICORP)
 - EC (Conquest, Qgates, Scala)
 - CNRS
 - UMPC
 - IUF
 - CdF